# MAGNETIC FLUX DETECTING OF BEARINGLESS INDUCTION MOTORS WITH SEARCH COILS 

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#### Abstract

The magnetic force of bearingless induction motors is generated with two kinds of magnetic fields interaction. The levitated mechanism determines the strong coupling relation between the air-gap flux of drive winding and that of levitation winding. Then it is very important to estimate or detect the distribution of magnetic flux. In this paper, as different positions of search coils are concerned, three methods of detecting the air gap flux density are analyzed in detail: evenly distributed angle position method, typical angle position method and analytic method. Moreover, the principle of the levitated force production is introduced. The control system configuration is proposed. At last, parameters of the motor prototype are given. The experimental results verify the validity of the proposed methods.


## INTRODUCTION

The interest in ultra-speed motors and generators in industry applications is increasing recently. Bearingless motors, which combine magnetic levitation bearings with the drive winding, have the advantage of no abrasion, no lubrication and no particle generation caused by bearings. Furthermore, in comparison with conventional magnetic bearing motors, the axial length of the rotor is decreased, which results in increased frequency of the mechanical natural vibration, and the efficiency can be increased because the magnetic flux of the motor can partly be used for levitation.

In recent years, a universal controller for bearingless induction motors has been proposed based on field oriented vector, such as air-gap flux orientation and rotor flux orientation ${ }^{[1] \sim[6]}$. However, there remain some limitations in the control algorithm:

1) Normal speed sensors are very expensive and none meets the measurement precision under the ultra-speed running. For example, the maximum speed of photoelectric encoder is lower than $10,000 \mathrm{r} / \mathrm{min}$.
2) There is great difficulty in realizing the adaptive
control or sensorless control, and restriction of the inherent maximal pullout torque because of complexity of the control computation. ${ }^{[7][8]}$
3) The control system requires an information link between the drive controller and the levitation one. Therefore, the low cost general-purpose inverter sold in the market cannot be used in this situation.

For these reasons, the independent control of levitated subsystem has to be built. In addition, in order to control the radial force, it is very important to detect or estimate the distribution of the magnetic flux densities in the air-gap of bearingless induction motors.

In early 2000, the method of detecting the amplitude and orientation of the air-gap magnetic field has been proposed with search coils wound around stator teeth of bearingless induction motors. ${ }^{[9][10]}$ However, these flux components calculation based on search coil has not reported in detail.

In this paper, three calculation methods of detecting the air-gap flux are analyzed in detail: evenly distributed angle position method, typical angle position method and analytic method. In other words, the air-gap flux densities $B_{1}$ of the drive winding and $B_{2}$ of the levitation winding can be detected respectively with search coils. To the author's best knowledge, the analytic method has not yet been published in literature before.

## PRINCIPLE AND CONTROL SYSTEM

According to [1], by combining 3-phase drive winding with pole pair number $p_{1}$ and 3-phase levitation winding with pole pair number $p_{2}=p_{1} \pm 1$ and keeping their angular velocity of revolving magnetic field: $\omega_{1}=\omega_{2}$, Maxwell force is generated for radial force in the induction motor.

Fig. 1 shows MMFs of two sets of windings of the bearingless induction motor. If $p_{2}=p_{1}+1$, the rotor angular velocity $\omega_{r}$ controlled by drive winding is


Fig. 1 MMFs of drive winding and levitation winding
larger than the synchronous angular velocity $\omega_{2}$, then the levitation winding operates as a generator. The stator's MMF $\vec{F}_{2 s}$ lags behind the air-gap's MMF $\vec{F}_{2}$. On the contrary, if $p_{2}=p_{1}-1$, the levitation winding operates as a motor. Then there is a phase delay of the rotor's MMF $\vec{F}_{2 r}$ to the air-gap's MMF $\vec{F}_{2}$. When the axes of A-phase of two windings overlap each other, their stators' MMFs are unanimous in the direction.

According to the general theory of motor control, if the drive winding is executed by sinusoidal current, its stator MMF $f_{1 s}(\varphi, t)$ can be written as

$$
\begin{equation*}
f_{1 s}(\varphi, t)=F_{1 s} \cos \left(\omega_{1} t-p_{1} \varphi\right) \tag{1}
\end{equation*}
$$

where $F_{1 \mathrm{~s}}$ is the peak value of the stator MMF, $\varphi$ is the space position angle.

The air-gap flux density produced by the drive winding is given by

$$
\begin{equation*}
B_{1}(\varphi, t)=B_{1 m} \cos \left(\omega_{1} t-p_{1} \varphi-\mu\right) \tag{2}
\end{equation*}
$$

where $B_{1 m}$ is the peak value of the air-gap flux density, $\mu$ is the phase angle from $\vec{F}_{1}$ to $\vec{F}_{1 s}$.

Similarly, the air-gap flux density of the levitation winding can be written as

$$
\begin{equation*}
B_{2}(\varphi, t)=B_{2 m} \cos \left(\omega_{2} t-p_{2} \varphi-\lambda\right) \tag{3}
\end{equation*}
$$

where $B_{2 m}$ is the peak value of the air-gap flux density; $\lambda$ is the phase angle from $\vec{F}_{2}$ to $\vec{F}_{2 s}$.

Then the total flux density in the air-gap is shown as (4), the summation of (2) and (3)

$$
\begin{align*}
B(\varphi, t) & =B_{1 m} \cos \left(\omega_{1} t-p_{1} \varphi-\mu\right) \\
& +B_{2 m} \cos \left(\omega_{2} t-p_{2} \varphi-\lambda\right) \tag{4}
\end{align*}
$$

Substituting equation (4) into the Maxwell Forces in $\alpha$ - and $\beta$-direction as follows

$$
\left.\begin{array}{l}
F_{\alpha}=\int_{0}^{2 \pi} \frac{l R}{2 \mu_{0}} B^{2}(\varphi, t) \cos \varphi d \varphi  \tag{5}\\
F_{\beta}=\int_{0}^{2 \pi} \frac{l R}{2 \mu_{0}} B^{2}(\varphi, t) \sin \varphi d \varphi
\end{array}\right\}
$$

where $l$ is the active length of the rotor, $R$ is the rotor radius, and $\mu_{0}$ is the vacuum or air permeability.

For the rotor current of the levitation winding is neglected, the air-gap flux density is expressed as

$$
\begin{equation*}
B_{2} \approx \frac{p_{2} L_{2 m}}{2 l R W_{2}} i_{2} \tag{6}
\end{equation*}
$$



Fig. 2 An independent control system of the levitation winding of bearingless induction motors
where $L_{2 m}$ is the mutual inductance of the levitation winding. $W_{2}$ is the number of turns for the levitation winding.

In this paper, $p_{1}=1, p_{2}=2$, then equation(5) can be simplified by equation (6) as

$$
\left[\begin{array}{c}
F_{\alpha}  \tag{7}\\
F_{\beta}
\end{array}\right]=\sqrt{\frac{1}{24}} \frac{\pi L_{2 m}}{W_{2} \mu_{0}}\left[\begin{array}{cc}
B_{1 \alpha} & B_{1 \beta} \\
-B_{1 \beta} & B_{1 \alpha}
\end{array}\right]\left[\begin{array}{c}
i_{2 \alpha} \\
i_{2 \beta}
\end{array}\right]
$$

Fig. 2 shows a block diagram of the proposed bearingless induction motor, the air-gap flux density components $B_{1 \alpha}$ and $B_{1 \beta}$ of the drive winding are calculated by detection block. In this situation, the control system of the levitation winding can be controlled independently, while the drive winding can be driven by general-purpose inverter.

## FLUX DETECTION

In the bearingless induction motor, two sets of 3 -phase windings are wound in the same stator slots: conventional $2 p_{1}$ poles drive winding and additional $2 p_{2}=2\left(p_{1} \pm 1\right)$ poles levitation winding. To detect the flux densities in the air-gap, search coils are wound around stator teeth at $\theta=2 \pi(j-1) / K, j=1,2, \cdots, K$. If $p_{1}$ is even, then $K=4 p_{1}$; otherwise $K=4 p_{2}$. The flux densities are detected by integrating induced voltage of search coils

For example, the two-pole ( $p_{1}=1$ ) driver winding and the four-pole ( $p_{2}=2$ ) levitation winding are wound in the stator of a typical 24 -slot induction motor in Fig.3, then search coils ( $K=4 p_{2}=8$ ) are wound around stator teeth at $\theta=0^{\circ}, 45^{\circ}, 90^{\circ}, 135^{\circ}, 180^{\circ}$, $225^{\circ}, 270^{\circ}$ and $315^{\circ}$.

The analysis below is based on the assumptions that the flux distribution in the air-gap confronted the stator tooth is uniform and the fringing flux is negligible.


Fig. 3 Stator winding arrangement in a motor prototype

When tooth flux density $B_{j}$ is detected by using search coil, search coil terminal voltage $V_{s c j}$ is given as

$$
\begin{equation*}
V_{s c j}=N_{t} S_{t} \frac{d}{d t} B_{j} \tag{8}
\end{equation*}
$$

where $N_{t}$ is the number of turns for a search coil; $S_{t}$ is the tooth cross section area; $j=1,2, \cdots, K$.

As a practical solution, a low-pass filter having a cut-off frequency $f_{c}$ and a DC gain $A$ are used as an incomplete integrator. Moreover, a high-pass filter having a time constant $T$ is needed for eliminating the integrator drift due to measurement-offset error. A transfer function between the terminal voltage $V_{s c j}$ and the actual magnetic flux density $B_{j}$ is given as

$$
\begin{equation*}
\frac{B_{j}}{V_{s c j}}=\frac{1}{N_{t} S_{t}} \frac{2 \pi f_{c} A}{s+2 \pi f_{c}} \frac{s}{s+\frac{1}{T}} \tag{9}
\end{equation*}
$$

where $s=d / d t$.
The magnetic flux distribution vectors $B_{1 \alpha}, B_{1 \beta}$ of the drive winding and $B_{2 \alpha}, B_{2 \beta}$ of the levitation winding are calculated from the detected magnetic flux densities $B_{j} . B_{j}$ can be calculated from the corresponding voltage $V_{s c j}$ according to equation(9).
A. Evenly Distributed Angle Position Method

In the negative ordinate Y-axis, the stationary angular position $\theta_{m}$ may be expressed as the corresponding $\theta_{n}$ in the positive Y-axis: $\theta_{m}=\theta_{n}+\pi$, where $m=n+K / 2, n=1,2, \cdots, K / 2$. Then

$$
\begin{align*}
& B_{m} \cos \left(p_{1} \theta_{m}\right)=B\left(\theta_{n}+\pi\right) \cos p_{1}\left(\theta_{n}+\pi\right) \\
& =B_{1 \alpha} \cos ^{2}\left(p_{1} \theta_{n}\right)+\frac{1}{2} B_{1 \beta} \sin \left(2 p_{1} \theta_{n}\right) \\
& -B_{2 \alpha} \cos \left(p_{2} \theta_{n}\right) \cos \left(p_{1} \theta_{n}\right)  \tag{10}\\
& -B_{2 \beta} \sin \left(p_{2} \theta_{n}\right) \cos \left(p_{1} \theta_{n}\right)
\end{align*}
$$

Substituting (10) into the summation as following

$$
\begin{align*}
& \sum_{j=1}^{K} B_{j} \cos \left(p_{1} \theta_{j}\right) \\
= & \sum_{n=1}^{K / 2} B_{n} \cos \left(p_{1} \theta_{n}\right)+\sum_{m=1+K / 2}^{K} B_{m} \cos \left(p_{1} \theta_{m}\right) \tag{11}
\end{align*}
$$

In this paper, $K=4 p_{2}, p_{2}$ is even number, then (11) can be simplified and the magnetic flux distribution vector $B_{1 \alpha}$ of the drive winding is written as

$$
\begin{equation*}
B_{1 \alpha}=\frac{2}{K} \sum_{j=1}^{K}\left[B_{j} \cos \left(p_{1} \theta_{j}\right)\right] \tag{12}
\end{equation*}
$$

In the same way, the other distribution vectors $B_{1 \beta}, B_{2 \alpha}$ and $B_{2 \beta}$ can be given as

$$
\begin{align*}
B_{1 \beta} & =\frac{2}{K} \sum_{j=1}^{K}\left[B_{j} \sin \left(p_{1} \theta_{j}\right)\right]  \tag{13}\\
B_{2 \alpha} & =\frac{2}{K} \sum_{j=1}^{K}\left[B_{j} \cos \left(p_{2} \theta_{j}\right)\right]  \tag{14}\\
B_{2 \beta} & =\frac{2}{K} \sum_{j=1}^{K}\left[B_{j} \sin \left(p_{2} \theta_{j}\right)\right] \tag{15}
\end{align*}
$$

By the evenly distributed angle position method, $B_{1 \alpha}, B_{1 \beta}, B_{2 \alpha}$ and $B_{2 \beta}$ are obtained by calculating flux densities of all the search coils. Therefore, the harmonic effect can be suppressed in certain degree and the precision in calculation is improved. However, the algorithm is comparatively complicated and difficult to realize in digital control. Furthermore, the method is ineffective if there is something wrong with coils at any angle position.
B. Typical Angle Position Method

Apparently, the flux density in a stator tooth positioned at $\theta$ is obtained based on

$$
\begin{align*}
B_{j} & =B_{1 \alpha} \cos \left(p_{1} \theta_{j}\right)+B_{1 \beta} \sin \left(p_{1} \theta_{j}\right) \\
& +B_{2 \alpha} \cos \left(p_{2} \theta_{j}\right)+B_{2 \beta} \sin \left(p_{2} \theta_{j}\right) \tag{16}
\end{align*}
$$

Then the air-gap flux densities toward the stator teeth at $0, \frac{\pi}{2}, \pi, \frac{3 \pi}{2}, B\left(\frac{\pi}{2 p_{2}}\right)$ and $B\left(\frac{\pi}{2 p_{2}}+\pi\right)$
are written as

$$
\left.\begin{array}{l}
B(0)=B_{1 \alpha}+B_{2 \alpha} \\
B(\pi)=-B_{1 \alpha}+B_{2 \alpha} \\
B\left(\frac{\pi}{2}\right)=(-1)^{\frac{p_{1}-1}{2}} B_{1 \beta}+(-1)^{\frac{p_{2}}{2}} B_{2 \alpha} \\
B\left(\frac{3 \pi}{2}\right)=(-1)^{\frac{p_{1}+1}{2}} B_{1 \beta}+(-1)^{\frac{p_{2}}{2}} B_{2 \alpha}  \tag{17}\\
B\left(\frac{\pi}{2 p_{2}}\right)=B_{1 \alpha} \cos \left(\frac{p_{1} \pi}{2 p_{2}}\right)+B_{1 \beta} \sin \left(\frac{p_{1} \pi}{2 p_{2}}\right)+B_{2 \beta} \\
B\left(\frac{\pi}{2 p_{2}}+\pi\right)=B_{2 \beta}-B_{1 \alpha} \cos \left(\frac{p_{1} \pi}{2 p_{2}}\right)-B_{1 \beta} \sin \left(\frac{p_{1} \pi}{2 p_{2}}\right)
\end{array}\right\}
$$

Sub-equations of (17) are added or subtracted, the magnetic flux distribution vectors $B_{1}$ of drive winding and $B_{2}$ of levitation winding in the perpendicular $\alpha$ - and $\beta$-axes are given as

$$
\left.\begin{array}{l}
B_{1 \alpha}=\frac{B(0)-B(\pi)}{2} \\
B_{1 \beta}=(-1)^{\frac{p_{1}-1}{2}} \frac{B\left(\frac{\pi}{2}\right)-B\left(\frac{3 \pi}{2}\right)}{2} \\
B_{2 \alpha}=\frac{B(0)+B(\pi)}{2}  \tag{18}\\
B_{2 \beta}=\frac{B\left(\frac{\pi}{2 p_{2}}\right)+B\left(\frac{\pi}{2 p_{2}}+\pi\right)}{2}
\end{array}\right\}
$$

Typical angle position method makes it possible to obtain the air-gap flux density components of two windings from the search coils at some certain angle positions. Each expression is simply determined by flux densities of search coils at two certain angular positions. Despite its advantage of easy realization, the results may be easily influenced by single search coil.
C. Analytic Method

In Fig.3, all of the air-gap flux densities confronted the search coils are obtained from (16) as

$$
\left[\begin{array}{l}
B_{1}  \tag{19}\\
B_{2} \\
B_{3} \\
B_{4} \\
B_{5} \\
B_{6} \\
B_{7} \\
B_{8}
\end{array}\right]=\left[\begin{array}{cccc}
1 & 0 & 1 & 0 \\
\frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} & 0 & 1 \\
0 & 1 & -1 & 0 \\
-\frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} & 0 & -1 \\
-1 & 0 & 1 & 0 \\
-\frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} & 0 & 1 \\
0 & -1 & -1 & 0 \\
\frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} & 0 & -1
\end{array}\right]\left[\begin{array}{l}
B_{1 \alpha} \\
B_{1 \beta} \\
B_{2 \alpha} \\
B_{2 \beta}
\end{array}\right]
$$

In the equations above, a linear equation with the variables of $B_{1} \sim B_{8}$ established by eliminating the magnetic flux $B_{1 \alpha}, B_{1 \beta}, B_{2 \alpha}$ and $B_{2 \beta}$ is as follows

$$
A\left[\begin{array}{llllllll}
B_{1} & B_{2} & B_{3} & B_{4} & B_{5} & B_{6} & B_{7} & B_{8} \tag{20}
\end{array}\right]^{\prime}=0
$$

where $A$ is a $4 \times 8$ matrix,

$$
A=\left[\begin{array}{cccccccc}
-1 & \frac{\sqrt{2}}{2} & 0 & -\frac{\sqrt{2}}{2} & 1 & -\frac{\sqrt{2}}{2} & 0 & \frac{\sqrt{2}}{2} \\
0 & \frac{\sqrt{2}}{2} & -1 & \frac{\sqrt{2}}{2} & 0 & -\frac{\sqrt{2}}{2} & 1 & -\frac{\sqrt{2}}{2} \\
1 & 0 & 1 & 0 & 1 & 0 & 1 & 0 \\
0 & 1 & 0 & 1 & 0 & 1 & 0 & 1
\end{array}\right]
$$

It is obvious that the rank of the matrix $A$ is four, $r(A)=4$. In other words, there are four $(K-r(A)=4)$ maximal linearly independent vectors (free vectors) in the fundamental solutions of equations, then the other vectors can be substituted by them. In addition, $B_{1 \alpha}, B_{1 \beta}, B_{2 \alpha}$ and $B_{2 \beta}$ can be calculated according to these free vectors. For example, the search coil flux
densities $B_{2}, B_{4}, B_{6}$ and $B_{7}$ at $\theta$ of $45^{\circ}, 135^{\circ}, 225^{\circ}$, and $270^{\circ}$ are maximal linearly independent vectors, then $B_{1}, B_{3}, B_{5}$ and $B_{8}$ can be written as

$$
\left[\begin{array}{l}
B_{1}  \tag{21}\\
B_{3} \\
B_{5} \\
B_{8}
\end{array}\right]=\left[\begin{array}{cccc}
-\frac{\sqrt{2}}{2} & -\sqrt{2} & -\frac{\sqrt{2}}{2} & -1 \\
\sqrt{2} & \sqrt{2} & 0 & 1 \\
-\frac{\sqrt{2}}{2} & 0 & \frac{\sqrt{2}}{2} & -1 \\
-1 & -1 & -1 & 0
\end{array}\right]\left[\begin{array}{l}
B_{2} \\
B_{4} \\
B_{6} \\
B_{7}
\end{array}\right]
$$

Substituting equation (21) into (19), yields

$$
\left[\begin{array}{l}
B_{1 \alpha}  \tag{22}\\
B_{1 \beta} \\
B_{2 \alpha} \\
B_{2 \beta}
\end{array}\right]=\left[\begin{array}{cccc}
0 & -\frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} & 0 \\
\frac{\sqrt{2}}{2} & \frac{\sqrt{2}}{2} & 0 & 0 \\
-\frac{\sqrt{2}}{2} & -\frac{\sqrt{2}}{2} & 0 & -1 \\
\frac{1}{2} & 0 & \frac{1}{2} & 0
\end{array}\right]\left[\begin{array}{l}
B_{2} \\
B_{4} \\
B_{6} \\
B_{7}
\end{array}\right]
$$

It is proved that ten combinations ( $B_{1}, B_{2}, B_{4}, B_{5}$ ), $\left(B_{1}, B_{2}, B_{6}, B_{7}\right),\left(B_{1}, B_{2}, B_{4}, B_{8}\right),\left(B_{1}, B_{3}, B_{5}, B_{7}\right),\left(B_{1}, B_{5}\right.$, $\left.B_{6}, B_{8}\right),\left(B_{2}, B_{3}, B_{5}, B_{6}\right),\left(B_{2}, B_{3}, B_{7}, B_{8}\right),\left(B_{2}, B_{4}, B_{6}, B_{8}\right)$, ( $\left.B_{3}, B_{4}, B_{6}, B_{7}\right),\left(B_{4}, B_{5}, B_{7}, B_{8}\right)$ are linearly dependent vectors, then there are totally $\mathrm{C}_{8}^{4}-10=60$ available combinations on maximal linearly independent vectors.

Analytic method takes full advantage of relations between the air-gap flux density and all flux densities of search coils. By using this method, the air-gap flux density components $B_{1 \alpha}, B_{1 \beta}, B_{2 \alpha}$ and $B_{2 \beta}$ can be calculated from flux densities of search coils corresponding to the maximal linearly independent vectors, and the number of necessary search coils is minimized to the amount of the maximal linearly independent vectors.

In fact, both evenly distributed angle position method and typical angle position method are special forms of analytic method. They are the special solutions composed of certain maximal linearly independent vectors and other vectors in the analytic method.

## EXPERIMENTAL RESULTS

In the paper, a bearingless induction motor prototype is designed. The parameters are as follows:

Active length of the rotor $l=9.5 \mathrm{~cm}$; rotor radius $R=14.5 \mathrm{~cm}$; weight of the rotor $m_{r}=1 \mathrm{~kg}$; clearance of the touch down bearing $\delta=200 \mu \mathrm{~m}$; number of turns for the drive winding $W_{1}=120$; number of turns for the levitation winding $W_{2}=32$; mutual inductance of the levitation winding $L_{2 m}=0.00435 \mathrm{H}$; cut-off frequency $f_{c} \approx 5 \mathrm{H}_{\mathrm{Z}}$; DC gain $A=0.414 \times 10^{-3}$; time constant $T=0.064 \mathrm{~s}$.

(a) Evenly distributed angle position

(c) Analytic method

(b) Typical angle position

(d) With the compensator

Fig. 4 Waveforms of the air-gap flux densities of the drive winding


Fig. 5 Waveforms of the independent levitated subsystem

To verify the validity of the methods proposed, the drive winding is controlled by the feasible air-gap flux orientation controller in [5][6]. Fig. 4 (a)-(c) show the commands flux densities $B_{1 \alpha}^{*}$ and $B_{1 \beta}^{*}$ of the drive winding, as well as the air-gap flux density components $B_{1 \alpha}$ and $B_{1 \beta}$ calculated according to flux densities of search coils. There is a phase delay between calculated results and commands caused by the incomplete integrator, the filter and $\mathrm{D} / \mathrm{A}$ sampling delay, etc. Then a phase lead-lag compensator is designed, and its effectiveness is shown in Fig. 4 (d).

In addition, the drive winding is driven directly by normal 3-phase alternating-current supply. Therefore, the independent control of levitated subsystem with search coils has been successfully built according to Fig.2. In Fig.5, the air-gap flux densities $B_{1}$ of motor winding and $B_{2}$ of the levitation winding are calculated by the analytic method. The waveforms are shown with the radial displacements of the rotor shaft at the rated speed of $2700 \mathrm{r} / \mathrm{min}$.

## CONCLUSION

The independent control between the levitation
and drive subsystem is an effective way for bearingless induction motors in ultra-speed operation. In this paper, three methods of detecting the flux density vectors $B_{1}$ of the drive winding and $B_{2}$ of the levitation winding with search coils are proposed, as well as the control system configuration. In that case, not only the ultra-speed operation of the bearingless induction motor becomes possible, but also the effect on levitation controlling, which caused by the drive winding's parameters change or error, is avoided. The experimental results verify the validity of the proposed methods.

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